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MEMORANDUM FOR PRS (In-House /Contractor Publication)

FROM: PROI (TI) (STINFO)

21 October 1999

SUBJECT: Authorization for Release of Technical Information, Control Number: AFRL-PR-ED-TP-1999-0196 Smith, C.W.; Gloss, K.T., Liu, C.T, "Test Geometries for Bondline Cracked Photoelastic Models; Preliminary Results" (VuGraphs)

ASME 1999 Mechanical Engineering Congress and Exposition

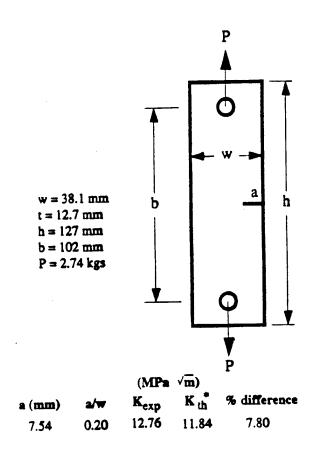
(Statement A)

TEST GEOMETRIES FOR BONDLINE CRACKED PHOTOELASTIC MODELS: PRELIMINARY RESULTS

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* Srawley and Brown, 1967

Fig. 1 Single Edge Crack Results for Artificial Cracks

(Mode I Algorithm)

Beginning with the Griffith-Irwin Equations, we may write, for Mode I, for the homogeneous case,

$$\sigma_{ij} = \frac{K_1}{(2\pi r)^{\frac{1}{2}}} f_{ij}(\theta) + \sigma_{ij}^{\diamond} \qquad (i.j. = n, z) \quad (1)$$

where:

 σ_{ij} are components of stress,

 K_1 is SIF,

 r, θ are measured from crack tip(Fig. A-1),

 σ_{ij}° are nonsingular stress components.

Then, along $\theta = \pi/2$ the direction of greatest local fringe spreading, after truncating σ_{ij}

$$(\tau_{nz})_{\text{max}} = \frac{K_1}{(8\pi r)^{\frac{1}{2}}} + \tau^{\circ} = \frac{K_{AP}}{(8\pi r)^{\frac{1}{2}}}$$
 (2)

where $\tau^{\circ} = f(\sigma_{ij}^{\circ})$ and is constant over the data range, $K_{AP} =$ apparent SIF, $(\tau_{nz})_{mas} =$ maximum shear stress in nz plane

$$\frac{K_{AP}}{\partial (\pi a)^{\frac{1}{2}}} = \frac{K_1}{\partial (\pi a)^{\frac{1}{2}}} + \frac{\sqrt{8}\tau^o}{\partial \left(\frac{r}{a}\right)^{\frac{1}{2}}}$$
(3)

where (Fig. A-1) a = crack length, and $\bar{\sigma} = \text{remote}$ normal stress

i.e.
$$\frac{K_{AP}}{\bar{\sigma}(\pi a)^{\frac{1}{2}}}$$
 vs. $\sqrt{\frac{r}{a}}$ is linear.

Since from the Stress-Optic Law:

$$(\tau_{nz})_{\max} = \frac{nf}{2t}$$
 where
 $n = \text{stress fringe order}$
 $f = \text{material fringe value}$
 $t = \text{specimen thickness}$

and from Eq. 2

$$K_{AP} = \tau_{ns}^{max} (8\pi r)^{\frac{1}{2}} = \frac{nf}{2t} (8\pi r)^{\frac{1}{2}},$$

then K_{AP} (through a measure of n) and r becomes the measured quantity from the stress fringe pattern at different points in the pattern.

A typical plot of normalized K_{AP} vs. $\sqrt{r/a}$ for a cracked, bonded specimen is shown in Fig. A-2.



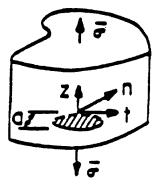


Fig. A-1 Mode I Notation

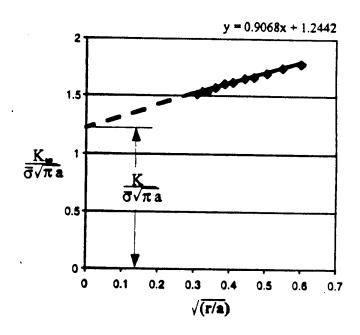


Fig. A-2: Determination of K₁ from Test Data for DS4.

The mixed mode algorithm was developed (see Fig. 12(a) and (b)) by requiring that

$$\lim_{\substack{m \to 0 \\ \Theta_m \to \Theta_m}} \left\{ (8\pi r_m)^{1/2} \frac{\delta(\tau)_{nz}^{\max}}{\delta \Theta} (K_1, K_2, r_m, \Theta_m, \tau_{ij}) \right\} = 0 \tag{4}$$

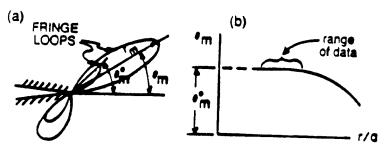


Fig. 12. (a) and (b). Determination of θ^{o}_{m} .

which leads to

$$\left(\frac{K_2}{K_1}\right)^2 - \frac{4}{3} \left(\frac{K_2}{K_1}\right) \cot 2\Theta_m^\circ - \frac{1}{3} = 0 - --.$$
 (5)

By measuring Θ°_{m} which is approximately in the direction of the applied load, K_{2}/K_{1} can be determined.

Then writing the stress optic law as

$$\tau_{nz}^{\max} = \frac{fn}{2t} = \frac{K_{AP}^{\bullet}}{(8\pi r)^{\frac{1}{2}}},$$

one may plot $K_{AP}^*/\overline{\sigma}(\pi a)^{1/2}$ vs $\sqrt{r/a}$ as before; locate a linear zone and extrapolate to r=0 to obtain K^* . Knowing, K^* , K_2/K_1 and Θ^* , walues of K_1 and K_2 may be determined since

$$K^* = \left[\left(K_1 \sin \Theta_m^{\circ} + 2K_2 \cos \Theta_m^{\circ} \right)^2 + \left(K_2 \sin \Theta_m^{\circ} \right)^2 \right]^{\frac{1}{2}} - --.$$
 (6)

Knowing K^* and Θ°_{m} , K_1 and K_2 can be determined from Eqs. (5) and (6). Details are found in Ref. [3].

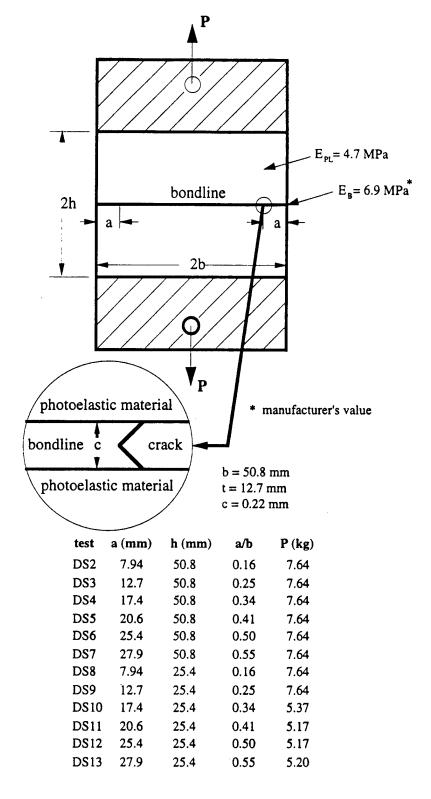


Fig. 2 Bonded Specimens with Double Edge Bondline Cracks

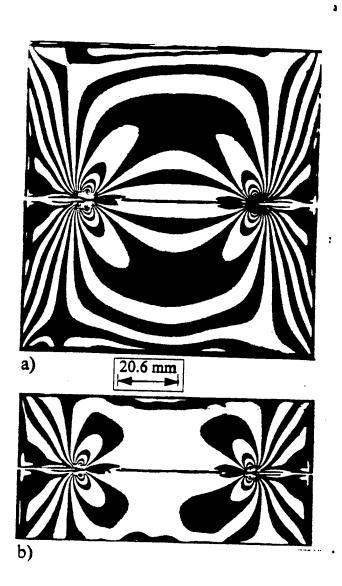


Fig. 3*: Global Stress Fringe Patterns for a) Square Specimen, b) Short Specimen.

[†]All fringe patterns have a bright background, (i.e. integral fringes are white, half fringes are black).

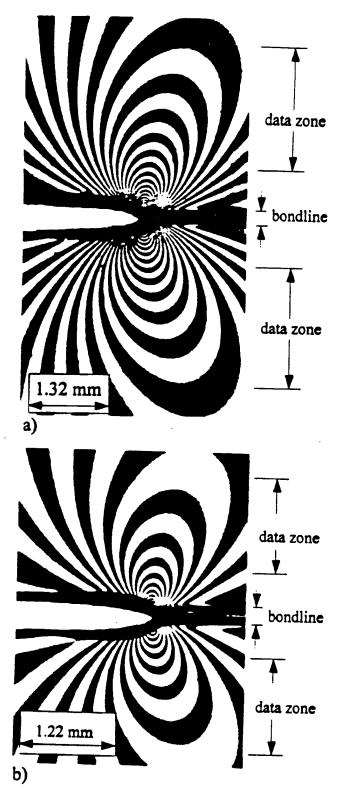


Fig. 4*: Local Stress Fringe Patterns for a) Square Specimen, b) Short Specimen.